

The Crystal and Molecular Structure of Sulfur Nitride, S₄N₄

BY BRAHAMA D. SHARMA* AND JERRY DONOHUE

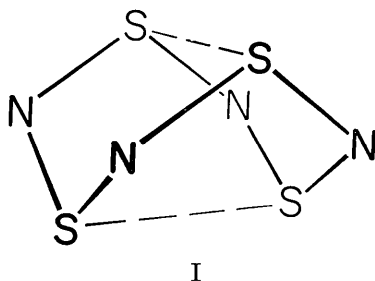
Department of Chemistry, University of Southern California, Los Angeles 7, California, U.S.A.

(Received 28 June 1962)

A three-dimensional refinement of crystalline S₄N₄ shows that molecule consists of a bisphenoid of sulfur atoms and a square of nitrogen atoms, in agreement with a previous electron diffraction investigation and a partial X-ray study. The inverse structure, *i.e.* a bisphenoid of nitrogen atoms and a square of sulfur atoms, which was supported by a different electron diffraction study and by spectroscopic investigations, is ruled out. The average S–N bond length is 1.616 ± 0.010 Å. The molecule may be described as a slightly elongated tetrahedron of sulfur atoms with nitrogen atoms added out from four of the edges.

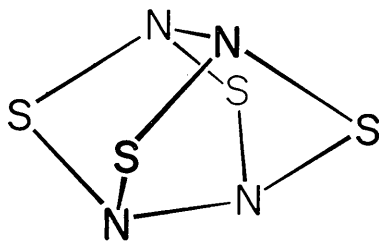
Introduction

Numerous structures were proposed for the S₄N₄ molecule before the advent of diffraction and related methods of structure determination. These all were listed and discussed by Lu & Donohue (1944) in their report of an electron diffraction study in which they favored a structure (I) which had not been suggested



I

previously. Just previously, Hassel & Viervoll (1943), also using the electron diffraction method, favored a different structure (II) which *prima facie* is more



II

reasonable than (I) because sulfur and nitrogen show their expected ligancies of two and three. For both structures the molecular symmetry is $\bar{4}2m$ (D_{2d}). The agreement between observed and calculated scattering curves was apparently satisfactory in both investigations, and since the spectra of interatomic distances

are quite different for the two structures, this conflict in the results is difficult to understand.

There have been three previous X-ray studies of S₄N₄. Jaeger & Zanstra (1931) reported that the crystals were orthorhombic with $a=8.87$, $b=8.47$, and $c=7.20$ (kX), and that the S₄N₄ molecule consisted of two interpenetrating bisphenoids of sulfur and nitrogen atoms with S–N = 1.2 Å. These results are questionable not only because of the remarkably short S–N bond distance but also because it had been previously shown by goniometry that crystalline S₄N₄ is monoclinic (Artini, 1907; Smith, 1911). In a later X-ray study, Buerger (1936) verified that the crystals were monoclinic, with $a=8.74$, $b=7.14$, $c=8.645$ (kX), $\beta=92^\circ 21'$, $Z=4$, space group $P2_1/n$ (C_{2h}^5). Finally, Clark (1952) verified the unit cell and space group reported by Buerger, and by a Patterson & Fourier analysis of the three prism zones obtained satisfactory agreement ($R=20\%$) assuming structure (I). He was apparently unaware of the work of Hassel & Viervoll.

Evidence favoring structure (II) was then presented by Lippincott & Tobin (1953), who studied the infrared and Raman spectra of both S₄N₄ and S₄N₄H₄. Their conclusion was based partly on indications of moderately high molecular symmetry together with the appearance of a Raman line at 888 cm^{-1} for S₄N₄, as compared with a line at 893 cm^{-1} from hydrazine corresponding to the N–N bond. They also concluded that the S₄N₄ and S₄N₄H₄ molecules were 'analogous in structure, except that in nitrogen hydrogen sulfide the two N–N bonds of nitrogen tetrasulfide have been broken to form N–H bonds.'

Both structures are nonpolar and therefore disagree with the dipole moment of 0.72 D reported by Phalnikar & Bhide (1939).

A recent evaluation of the parachor of S₄N₄ (Krylov, Patsuk & Podchainova, 1960) was said to be in good agreement with values calculated for both structures, and thus is of no aid in settling the controversy.

In the present paper we report the results of a three-dimensional X-ray study which show the S₄N₄ molecule to have structure (I).

* Present address: Gates and Crellin Laboratories of Chemistry, California Institute of Technology, Pasadena 4, California, U.S.A.

Experimental

A sample of S_4N_4 , prepared by the method described by Arnold, Hugill & Hutson (1936), was kindly supplied by Prof. N. Kharasch of this department, under whose directions our own preparations of some pure material had been discarded. Numerous attempts to obtain single crystals were unsuccessful, but an acicular twin, of approximately 0.1 mm radius, was finally obtained which gave a diffraction pattern which could be indexed unequivocally. Multiple film Weissenberg photographs ($Cu K\alpha$) about b , $k=0$ to 4, were taken, the intensities being estimated visually by means of an intensity strip. A total of 651 unique reflections was recorded. Correlation between the layers was effected with $(0kl)$ and $(hk0)$ precession photographs of the same crystal. After correction of the intensities for absorption with the powder rod formula, relative F values were calculated in the usual way. The final values of the interlayer scale factors were averages of the values obtained (a) with the precession photographs, (b) from application of Wilson's method to the individual layers, and (c) comparison of the F_o with those calculated using the parameters of Clark (1952). These last calculations, in addition to giving an isotropic B value of 2.67 \AA^2 , also suggested that Clark's structure was essentially correct, since for the layers $k=0$ to 4 the respective values of R were 24%, 28%, 22%, 21%, and 29%. The lattice constants $a=8.75$, $b=7.16$, $c=8.65 \text{ \AA}$, $\beta=92.5^\circ$ (Clark) have been adopted in the present study.

Refinement

The structure was refined by both least-squares and Fourier methods. The parameters of Clark were used as the starting point in a least-squares treatment in which the positional parameters and individual isotropic B 's were refined. The quantity minimized was $\sum w(F_o - F_c)^2$. The weighting scheme was that of Hughes (1941). The form factor curve for sulfur was that of Dawson (1961), and for nitrogen, that of Berghuis, Haanappel, Potters, Loopstra, MacGillavry & Veenendaal (1955). The resulting values were then subjected to further refinement by the use of three-dimensional difference maps (Leung, Marsh & Schomaker, 1957), from which anisotropic B_{ij} were evaluated. No spurious peaks were noted in the

difference maps. The final values of the thermal and positional parameters were obtained from successive full matrix least squares. The shifts in the last cycle were all less than 10^{-4} (positional parameters as fractions of cell edges, B 's as coefficients of h^2 , etc.). During the above calculations R (absent reflections omitted) dropped steadily from 23.2% to 9.8%. Observed and calculated structure factors are presented in Table 1. The final parameters are presented in Table 2. The respective average and maximum changes from the starting positional parameters are, for the sulfur atoms, 0.032 and 0.055 \AA , and, for the nitrogen atoms, 0.055 and 0.132 \AA .

Discussion

Molecular structure

The intramolecular distances and angles are shown in Table 3 and Fig. 1. The molecular symmetry is 4_2m within the limits of error. The average S-N bond distance is $1.616 \pm 0.010 \text{ \AA}$. The previously reported

Table 3. *Intramolecular distances and angles*

The standard errors in the distances have been multiplied by 10^3

Distance	(σ)	Angle	(σ)
S(1)-N(4)	1.631 \AA (13)	N(2)-S(1)-N(4)	103.3° (1.6°)
S(2)-N(4)	1.605 (19)	N(3)-S(3)-N(1)	104.9 (1.6)
S(2)-N(1)	1.621 (14)	N(1)-S(2)-N(4)	105.1 (1.5)
S(3)-N(1)	1.616 (17)	N(2)-S(4)-N(3)	104.6 (1.3)
S(3)-N(3)	1.596 (14)		
S(4)-N(3)	1.617 (22)	S(2)-N(1)-S(3)	112.3 (0.6)
S(4)-N(2)	1.610 (14)	S(4)-N(2)-S(1)	112.8 (0.7)
S(1)-N(2)	1.634 (18)	S(3)-N(3)-S(4)	113.6 (0.7)
S(1) \cdots S(3)	2.586 (8)	S(1)-N(4)-S(2)	112.6 (0.7)
S(2) \cdots S(4)	2.576 (8)		
S(1) \cdots S(2)	2.693 (8)		
S(1) \cdots S(4)	2.702 (10)		
S(3) \cdots S(4)	2.689 (8)		
S(2) \cdots S(3)	2.688 (6)		

values of $1.60 \pm 0.04 \text{ \AA}$ (Clark, 1952) and $1.62 \pm 0.02 \text{ \AA}$ (Lu & Donohue, 1944) are not significantly different from this value. All of these disagree with value 1.74 \AA reported by Hassel & Viervoll (1943). The intramolecular distances S(1) \cdots S(3) and S(2) \cdots S(4) are much shorter than would be expected on the basis of the van der Waals diameter of $ca 3.3 \text{ \AA}$, but are much longer than the S-S single bond distance of 2.05 \AA . The S-N bond distance corresponds to about

Table 2. *Final parameters and their standard errors*

All values have been multiplied by 10^4

Atom	$x(\sigma_x)$	$y(\sigma_y)$	$z(\sigma_z)$	$b_{11}(\sigma_{b_{11}})$	$b_{22}(\sigma_{b_{22}})$	$b_{33}(\sigma_{b_{33}})$	$b_{12}(\sigma_{b_{12}})$	$b_{13}(\sigma_{b_{13}})$	$b_{23}(\sigma_{b_{23}})$
S(1)	106(4)	9260(6)	2996(4)	118(4)	156(12)	135(5)	19(12)	-6(6)	-157(13)
S(2)	1500(3)	7088(5)	943(4)	66(3)	97(11)	128(4)	37(9)	17(5)	-18(10)
S(3)	-1488(3)	7951(6)	652(4)	66(3)	112(12)	133(4)	47(9)	-27(5)	0(11)
S(4)	-353(4)	5535(6)	2787(4)	126(5)	137(13)	135(5)	-10(12)	21(7)	116(12)
N(1)	70(12)	7587(17)	-238(12)	115(13)	116(30)	105(13)	112(35)	-20(20)	-79(35)
N(2)	-172(14)	7309(21)	3923(13)	156(18)	278(47)	116(14)	-11(48)	24(26)	10(46)
N(3)	-1761(11)	6075(19)	1597(14)	87(13)	91(32)	185(19)	13(34)	23(24)	85(41)
N(4)	1663(11)	8851(19)	2083(13)	83(11)	114(32)	156(16)	-12(33)	-33(21)	-141(38)

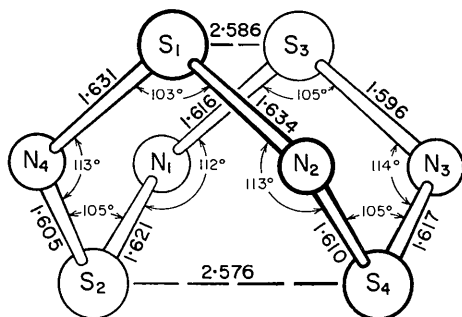
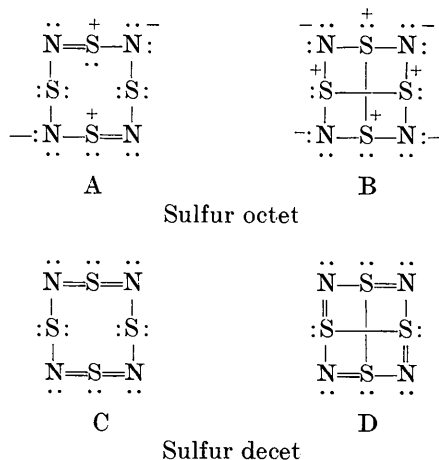


Fig. 1. Bond distances and angles.

one-third double bond character. The bond angles of 105° at a sulfur atom and 113° at a nitrogen atom, on the other hand, do not imply very much, if any, double bond character. This paradox scarcely seems resolved on consideration of what are probably the most important of the many possible resonance structures:



While these considerations allow a discussion of the results, they scarcely can be said to explain why the S₄N₄ molecule assumes the unusual structure that it does. It seems likely that the geometrical restraints

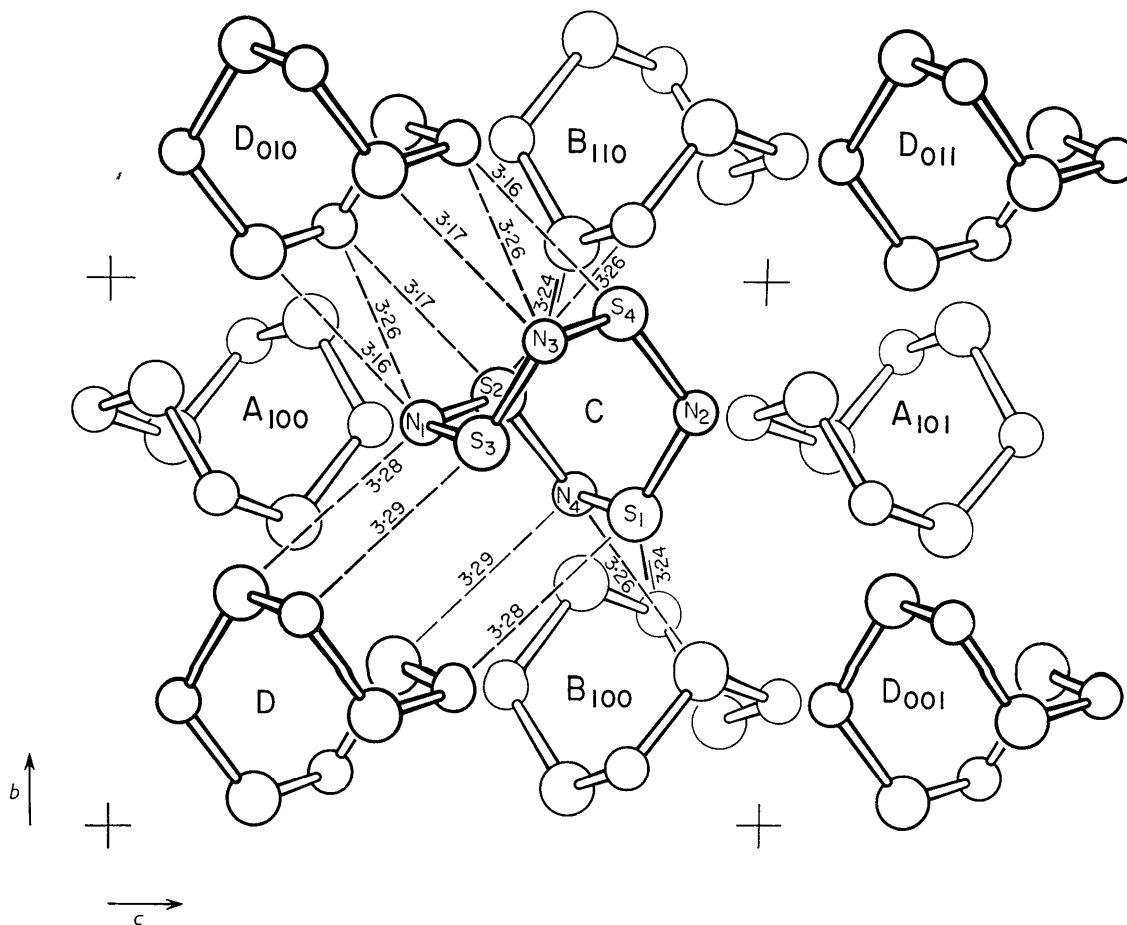


Fig. 2. The structure viewed down the *a* axis.

imposed by the molecular symmetry may well result in some strain, a conclusion not in disagreement with the rather unstable nature of the substance.

A somewhat different interpretation has been proposed by Lindquist (1958), who suggested, partly on the basis of the geometry of the $\begin{matrix} \text{N} \\ \text{N} \end{matrix} \text{S} \cdots \text{S}$ groupings (the eight different N-S \cdots S angles are all within 2° of 90°), that the S \cdots S distances of 2.58 Å corresponded to pure $p\sigma$ bonds, and that only structure B was important. While this explanation has its attractions, it does not appear to account for the short S-N bond length. Other molecules in which the S-N bond is formally single include sulfamic acid, $\text{H}_3\text{N}^+\text{SO}_3^-$ (Sass, 1960) and $(\text{SNH})_4$ (Sass & Donohue, 1958) in which the respective bond distances are 1.764 ± 0.020 Å and 1.674 ± 0.004 Å; the former is probably a pure single (sp^3) bond, and the latter then corresponds to some double ($p\pi$) bond character, a situation also suggested by the bond angles, as discussed by Sass & Donohue. The single-bond radius sum is 1.74 Å (Schomaker & Stevenson, 1941).

A reinterpretation of the spectroscopic data reported by Lippincott & Tobin in terms of structure (I), rather than structure (II) favored by them, is to be desired. The same suggestion applies to their spectral data for $(\text{SNH})_4$, since that molecule is a staggered ring with symmetry $4mm$, which is hardly 'analogous' to the $\bar{4}2m$ symmetry of the S_4N_4 molecule. Similarly, a redetermination of the dipole moment is desirable.

Molecular packing

The arrangement of the molecules in the crystal roughly approximates cubic close packing, as may be seen from Figs. 2 and 3. All intermolecular distances less than 3.30 Å are indicated in Fig. 2, and all those less than 3.50 Å listed in Table 4. These are all seen to be S \cdots N distances; the shortest S \cdots S distance is 3.64 Å, a value greater than the numerous van der Waals contacts of 3.50–3.53 Å observed in S_6 (Donohue, Caron & Goldish, 1961). It thus seems likely that in S_4N_4 the molecular packing is determined by intermolecular S \cdots N contacts. The shortest S \cdots N contacts of 3.16 Å are 0.2 Å shorter than the sum of the van der Waals radii (Pauling, 1960)

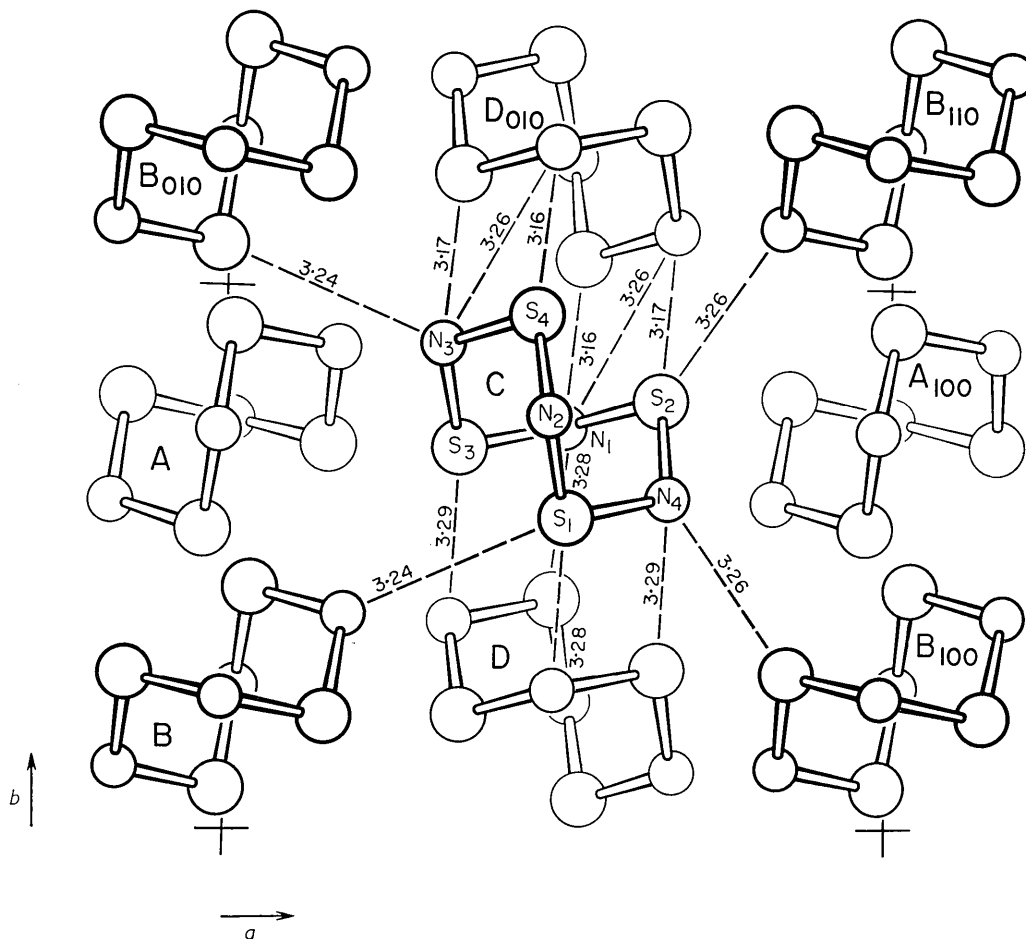


Fig. 3. The structure viewed down the c axis.

Table 4. Intermolecular distances

Molecular designations are $A(x, y, z)$, $B(-x, 1-y, 1-z)$, $C(\frac{1}{2}+x, \frac{3}{2}-y, \frac{1}{2}+z)$, $D(\frac{1}{2}-x, -\frac{1}{2}+y, \frac{1}{2}-z)$; subscripts give lattice translations from these

From atom in molecule C		To atom	in molecule	Distance	Equivalent distance		
					Atom in C	to atom	in molecule
S(1)		N(3)	B	3.235 Å	N(3)	S(1)	B ₀₁₀
S(1)		N(1)	D	3.283	S(1)	N(1)	D
S(2)		N(2)	A ₁₀₀	3.487	N(2)	S(2)	A ₀₀₁
S(2)		N(4)	B ₁₁₀	3.260	N(4)	S(2)	B ₁₀₀
S(2)		N(3)	D ₀₁₀	3.170	N(3)	S(2)	D ₀₁₀
S(3)		N(2)	A	3.498	N(2)	S(3)	A ₁₀₁
S(3)		N(1)	D	3.452	N(1)	S(3)	D
S(3)		N(4)	D	3.291	N(4)	S(3)	D
S(4)		N(4)	B ₁₁₀	3.444	N(4)	S(4)	B ₁₀₀
S(4)		N(1)	D ₀₁₀	3.157	N(1)	S(4)	D ₀₁₀
N(1)		N(3)	D ₀₁₀	3.255	N(3)	N(1)	D ₀₁₀
N(1)		N(4)	D	3.338	N(4)	N(1)	D
N(1)		N(1)	D	3.482	N(1)	N(1)	D ₀₁₀

Table 5. *R.M.S. amplitudes* $(\overline{u^2})^{\frac{1}{2}}$ and *direction cosines* q relative to a^*bc of the principal axes of the temperature factor ellipsoids

Atom	Axis i	$(\overline{u^2})^{\frac{1}{2}}$	q_i^1	q_i^2	q_i^3
S(1)	1	0.268	-0.196	-0.620	0.760
	2	0.212	-0.980	0.090	-0.179
	3	0.144	-0.043	0.780	0.625
S(2)	1	0.221	0.019	-0.114	0.993
	2	0.177	0.723	0.688	0.065
	3	0.138	0.691	-0.717	-0.096
S(3)	1	0.229	-0.292	-0.093	0.952
	2	0.185	-0.537	-0.808	-0.243
	3	0.138	-0.792	0.582	-0.186
S(4)	1	0.251	-0.013	-0.551	-0.835
	2	0.221	-0.993	0.104	-0.054
	3	0.152	-0.116	-0.828	0.548
N(1)	1	0.252	-0.696	-0.555	0.456
	2	0.190	-0.557	0.015	-0.831
	3	0.124	-0.454	0.832	0.319
N(2)	1	0.269	0.137	-0.990	-0.041
	2	0.246	0.977	0.128	0.171
	3	0.208	0.164	0.064	-0.985
N(3)	1	0.271	0.042	0.259	0.965
	2	0.183	0.996	0.067	-0.061
	3	0.141	0.080	-0.964	0.255
N(4)	1	0.269	0.187	0.449	-0.874
	2	0.182	0.934	-0.357	0.016
	3	0.127	0.305	0.819	0.486

of 1.5 for N and 1.85 for S. It is interesting that an anomalously short S...S distance (of 3.37 Å) also occurs in orthorhombic S₈ (Donohue *et al.*, 1961).

Thermal vibrations

The atom vibration ellipsoids corresponding to the thermal parameters of Table 2 are presented in Table 5. Unfortunately, there does not appear to be much correlation of these anisotropies with the symmetry axes of the molecule, so it seems quite possible that the B_{ij} either do not correspond to a physically significant situation or that the thermal motions arise in no small part from the nature of the packing and are therefore difficult to untangle from the motions of the individual molecules. In fact, calculation of the T, ω tensor (Cruckshank, 1956) gives a result which suggests that the interpretation of the thermal vibra-

tions of Table 5 in terms of rigid body motions is a rather dubious procedure. Accordingly, the evaluation of six B_{ij} per atom was, in the present case, probably merely a device for reducing the residual. The present example is doubtless not unique in this regard.

This work was supported by the National Science Foundation. The figures were prepared by Mrs Maryellin Reinecke. We wish to thank Dr Richard Marsh for helpful discussion and for advice and assistance with those calculations which were done on a Burroughs Datatron 205 at the California Institute of Technology. Other calculations, at the Western Data Processing Center, were done with the help of Dr A. Caron and programs written by Dr G. J. Palenik. The T, ω program was supplied by Prof. K. N. Trueblood.

References

- ARNOLD, M. H. M., HUGILL, J. A. C. & HUTSON, J. M. (1936). *J. Chem. Soc.* p. 1645.
 ARTINI, E. (1907). *Z. Kristallogr.* **42**, 68.
 BERGHUIS, J., HAANAPPEL, IJ. M., POTTERS, M., LOOPSTRA, B. O., MACGILLAVRY, C. H. & VEENENDAAL, A. L. (1955). *Acta Cryst.* **8**, 478.
 BUERGER, M. J. (1936). *Amer. Min.* **21**, 575.
 CLARK, D. (1952). *J. Chem. Soc.* p. 1615.
 CRUCKSHANK, D. W. J. (1956). *Acta Cryst.* **9**, 754.
 DAWSON, B. (1961). *Acta Cryst.* **14**, 1271.
 DONOHUE, J., CARON, A. & GOLDISH, E. (1961). *J. Amer. Chem. Soc.* **83**, 3748.
 HASSEL, O. & VIERVOLL, H. (1943). *Tidsskr. Kemi Bergv.* **3**, 7.
 HUGHES, E. W. (1941). *J. Amer. Chem. Soc.* **63**, 1737.
 JAEGER, F. M. & ZANSTRA, J. E. (1931). *Proc. Acad. Sci. Amst.* **34**, 782.
 KRYLOV, E. I., PATSUK, V. V. & PODCHAINOVA, V. N. (1960). *Izvest. Vysshikh Ucheb. Zavedenii, Khim. i Khim. Tekhnol.* **3**, 41.
 LEUNG, Y. C., MARSH, R. E. & SCHOMAKER, V. (1957). *Acta Cryst.* **10**, 650.
 LINDQUIST, I. (1958). *J. Inorg. Nucl. Chem.* **6**, 159.
 LIPPINCOTT, E. R. & TOBIN, M. C. (1953). *J. Chem. Phys.* **21**, 1559.

- LU, C.-S. & DONOHUE, J. (1944). *J. Amer. Chem. Soc.* **66**, 818.
- PAULING, L. (1960). *The Nature of the Chemical Bond*, 3rd ed., p. 260. Ithaca: Cornell University Press.
- PHALNIKAR, N. L. & BHIDE, B. V. (1939). *Curr. Sci.* **8**, 473.
- SASS, R. L. (1960). *Acta Cryst.* **13**, 320.
- SASS, R. L. & DONOHUE, J. (1958). *Acta Cryst.* **11**, 497.
- SCHOMAKER, V. & STEVENSON, D. P. (1941). *J. Amer. Chem. Soc.* **63**, 37.
- SMITH, G. F. H. (1911). *Miner. Mag.* **16**, 97.

Acta Cryst. (1963). **16**, 897

The Crystal and Molecular Structure of Taurine

BY H. H. SUTHERLAND* AND D. W. YOUNG

Carnegie Laboratory of Physics, Queen's College, Dundee, Scotland

(Received 21 August 1962)

The structure of taurine, $\text{NH}_2[\text{CH}_2]_2\text{HSO}_3$, has been determined by Patterson superposition and Fourier methods. The crystals are monoclinic with space group $P2_1/c$. The unit-cell dimensions are,

$$a = 5.30, b = 11.65, c = 7.94 \text{ \AA}, \beta = 94^\circ 17'.$$

The structure was refined by least-squares methods using three-dimensional data. The final R -factor, 0.11, was obtained using individual isotropic temperature factors for the heavier atoms and an over-all isotropic temperature factor for the hydrogen atoms. The bond lengths, bond angles and inter-molecular contacts were found to be normal.

Introduction

The X-ray investigation of several related compounds is at present being undertaken by the authors, to obtain structural information on the sulphonate and sulphuric radicals, taurine being the first of the series.

Taurine, $\text{NH}_2[\text{CH}_2]_2\text{HSO}_3$, occurs in the combined state in ox-gall and many other animal secretions. In the animal organism taurine is formed from cysteine by decarboxylation followed by oxidation. It can be obtained synthetically by the addition of sodium bisulphite to nitroethylene and the subsequent reduction of the nitroethane sulphonic acid. It forms colourless crystals, and melts and decomposes at about 328°C ; it is readily soluble in water but insoluble in alcohol and is quoted as forming tetragonal crystals (*Handbook of Chemistry and Physics*, 1961).

Experimental

A sample of taurine was recrystallized from water, the form of the resulting crystals depending on the rate of evaporation of the solution. With rapid evaporation, needle-shaped crystals were formed—the length of the needle being parallel to the a -axis—whereas with slow evaporation the crystals adopted a tabular appearance, the greatest length being along the a -axis. Both crystal types possessed identical space groups and unit-cell dimensions.

The space group and unit-cell dimensions were

obtained from Weissenberg films using $\text{Cu } K\alpha$ radiation; the systematic absences,

$$0k0 \text{ with } k=2n+1 \text{ and } h0l \text{ with } l=2n+1$$

uniquely determined the space group as $P2_1/c$; the unit-cell dimensions were:

$$a = 5.30 \pm 0.01, b = 11.65 \pm 0.01, \\ c = 7.94 \pm 0.01 \text{ \AA}; \beta = 94^\circ 17' \pm 5'.$$

The density measured at 20°C by the method of flotation, using a benzene-tetrabromoethane mixture, was $1.70 \pm 0.02 \text{ g.cm}^{-3}$. The calculated value, assuming four molecules to the unit cell, was 1.71 g.cm^{-3} .

The intensity data were measured visually from sets of multiple-film equi-inclination Weissenberg photographs (Robertson, 1943) taken with $\text{Cu } K\alpha$ radiation. 850 independent reflexions were measured from the $0kl$, $hk0$, $h0l$ and hkh zones and upper layers ($h=1$ to 5) taken with the a -axis as rotation axis. The intensities were corrected for Lorentz and polarization factors, but since the crystals were needle-shaped no attempt was made to apply absorption corrections. A correction for the variation in the area of the reflexions occurring in the upper-level Weissenberg photographs was made by the method of Phillips (1954). It was found that when the Phillips correction factor was applied to the extended reflexions, there was a difference in scale factor between the reflexions on both halves of the film, the scale factor appearing to vary with the Bragg angle, θ . An empirical correction curve was calculated for the contracted reflexions

* Present address: Physics Department, University of Hull, England.